

Poster # 1776

**Geometry meets feedback loops:
Shearing and turbulence self-regulation
in negative triangularity tokamaks**

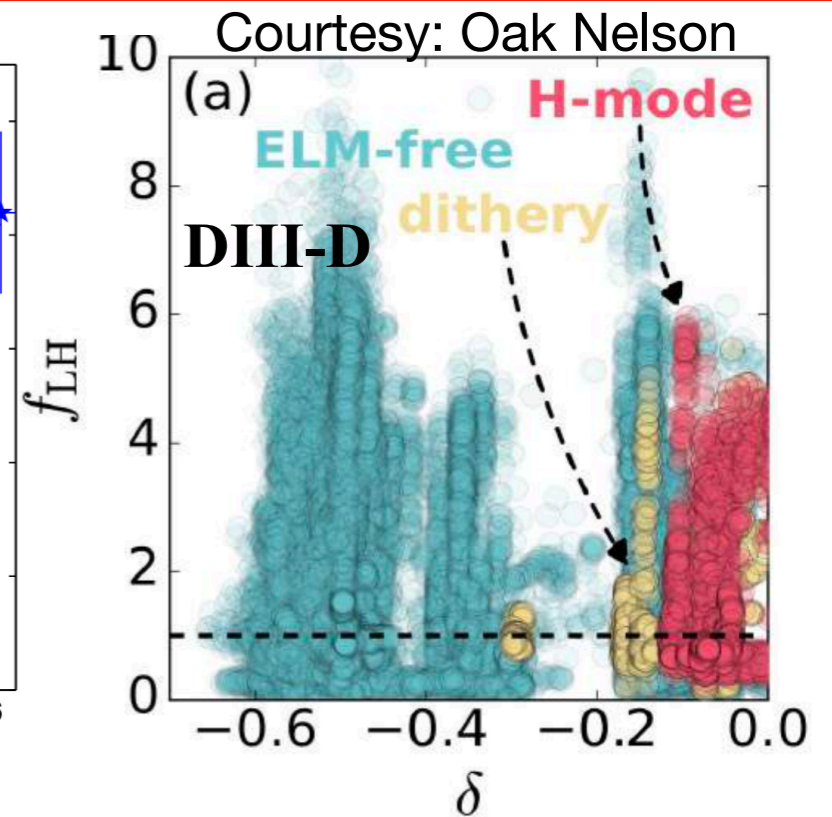
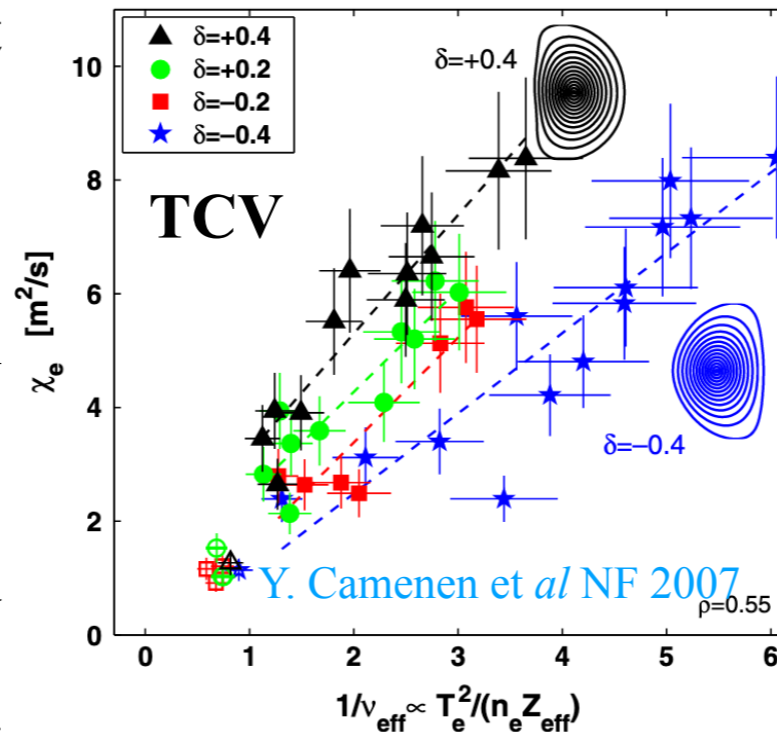
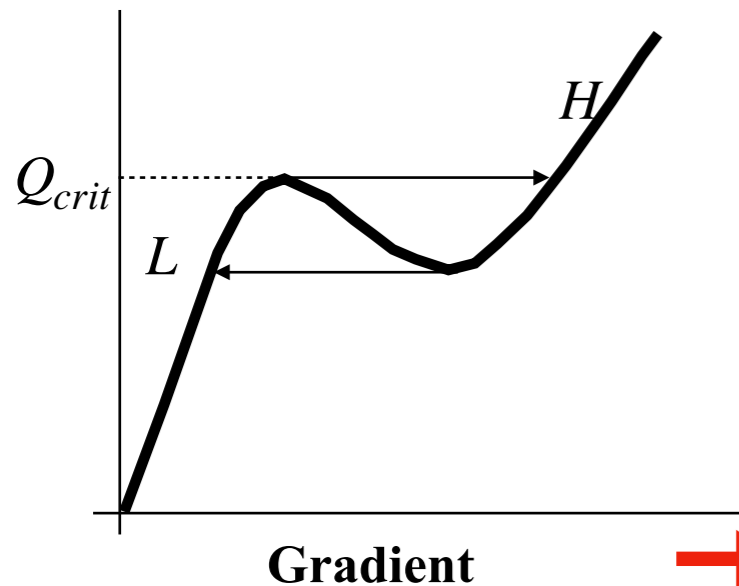
Rameswar Singh and P H Diamond
Department of Astronomy and Astrophysics, UCSD

29th Fusion Energy Conference Oct 16 - 21, 2023
London, UK

Acknowledgements: T. S. Hahm, O Nelson, L. Schmitz, K. Thome and “Neg. Triang. Cabalists ” DIII-D.
U.S. Department of Energy Award Number DE-FG02-04ER54738.

Effect of triangularity on confinement, fluctuations, and L-H transition

- TCV: Energy confinement time doubled, fluctuations reduced when $\delta \rightarrow -\delta$
- DIII-D: No H-mode transition for $\delta < \delta_{crit} \sim -0.18$
 - $P_{L \rightarrow H}$ diverges for $\delta < \delta_{crit}$
 - loss of access to 2nd stability region of $n = \infty$ ideal MHD ballooning modes for $\delta < \delta_{crit}$ [Saarelma *et al* PPCF 2021, Nelson *et al* NF 2022].



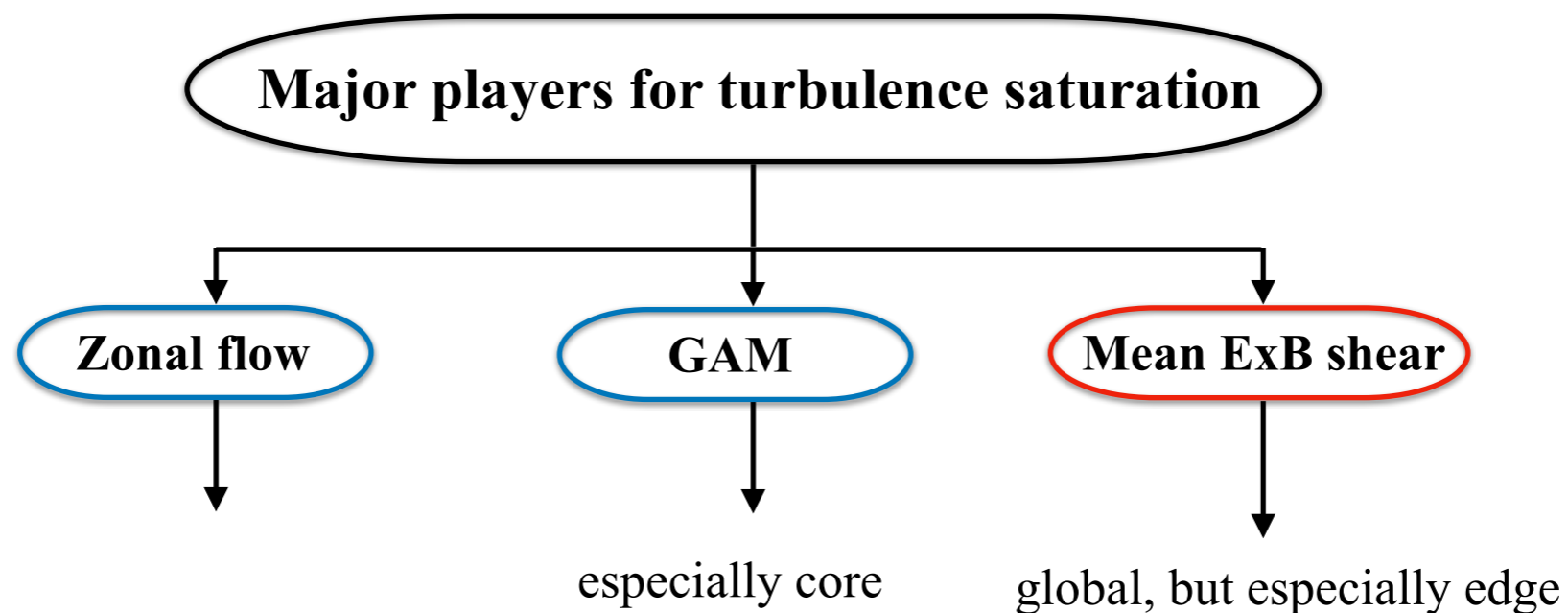
- Is H mode operation always in 2nd stability region?
 - Magnetic separatrix and finite edge current can cause coalescence of 1st and 2nd stable region. [Bishop NF 1986]
 - Many past examples of (PT) H mode operation in the 1st stable region.
 - H-mode persisted even after loss of 2nd stability. [Lao *et al* NF 1999, J R Ferron *et al* NF 2000]
- What happens to the E_r' induced transport bifurcation picture of L-H transition in NT?

→ Role of mean ExB shear in NT pedestal formation?

How to reconcile confinement improvement in NT L-mode with diverging

$$P_{th,L \rightarrow H} \text{ for } \delta < \delta_{crit} ?$$

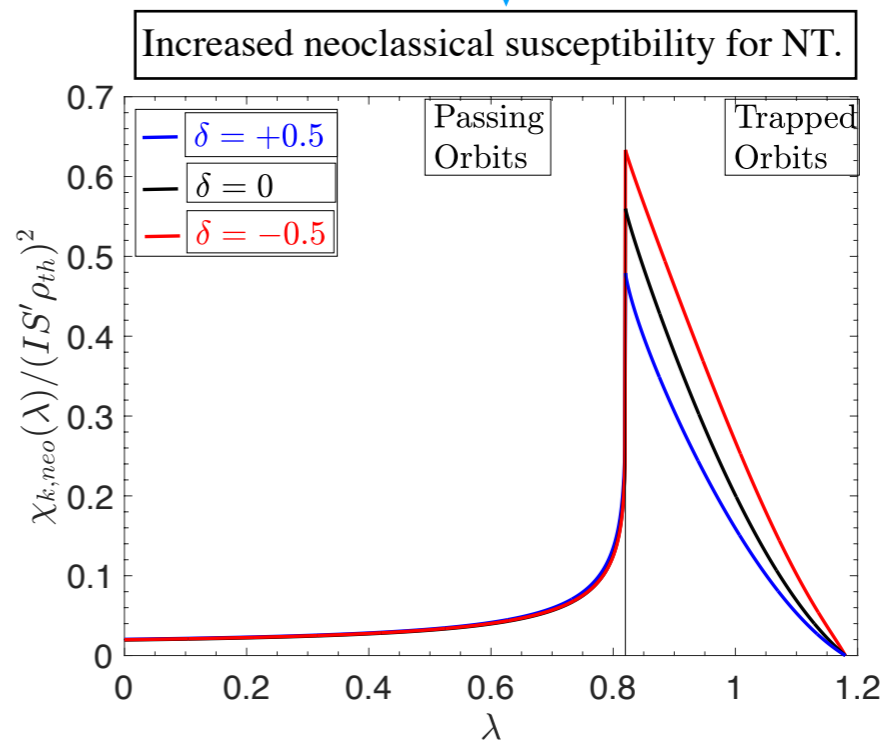
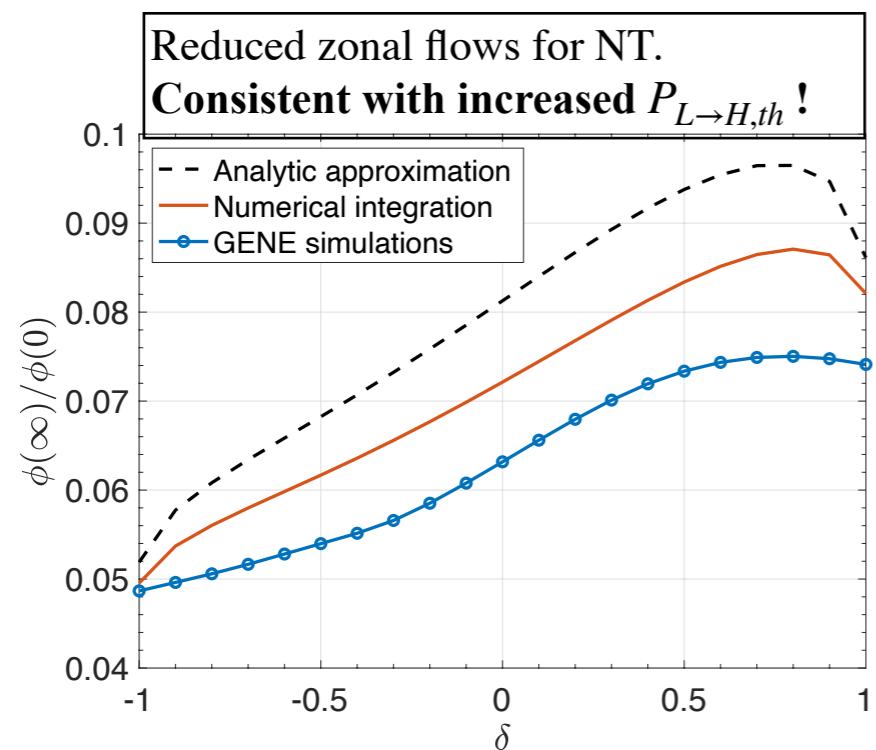
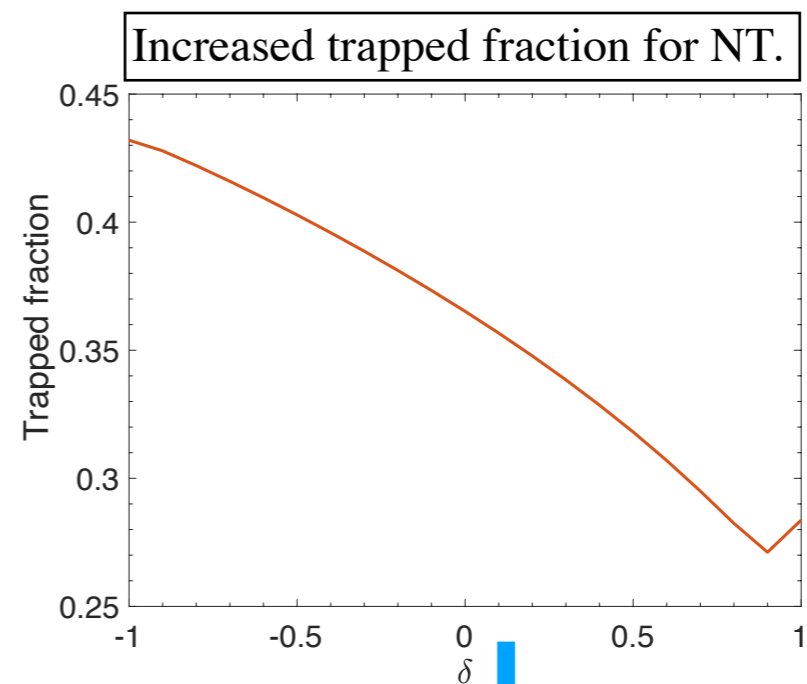
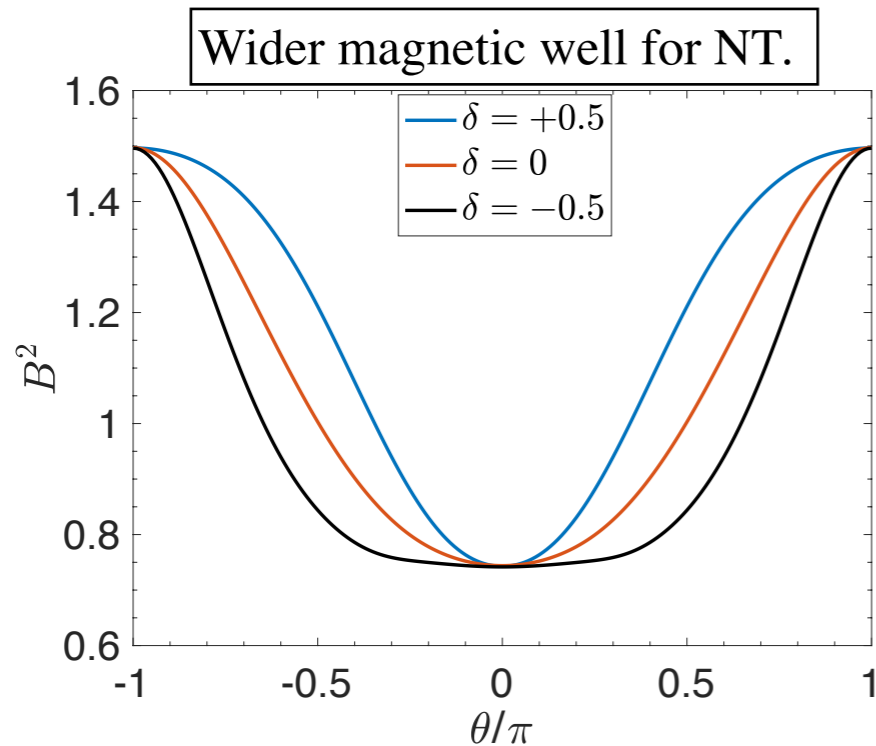
- Need think beyond linear stabilization of zoo of modes(TEM/ITG,...)!
- Understanding flux surface shaping effects on turbulence saturation mechanism is important.



- *Interplay of NT configuration with secondary modes feedback and shearing?*

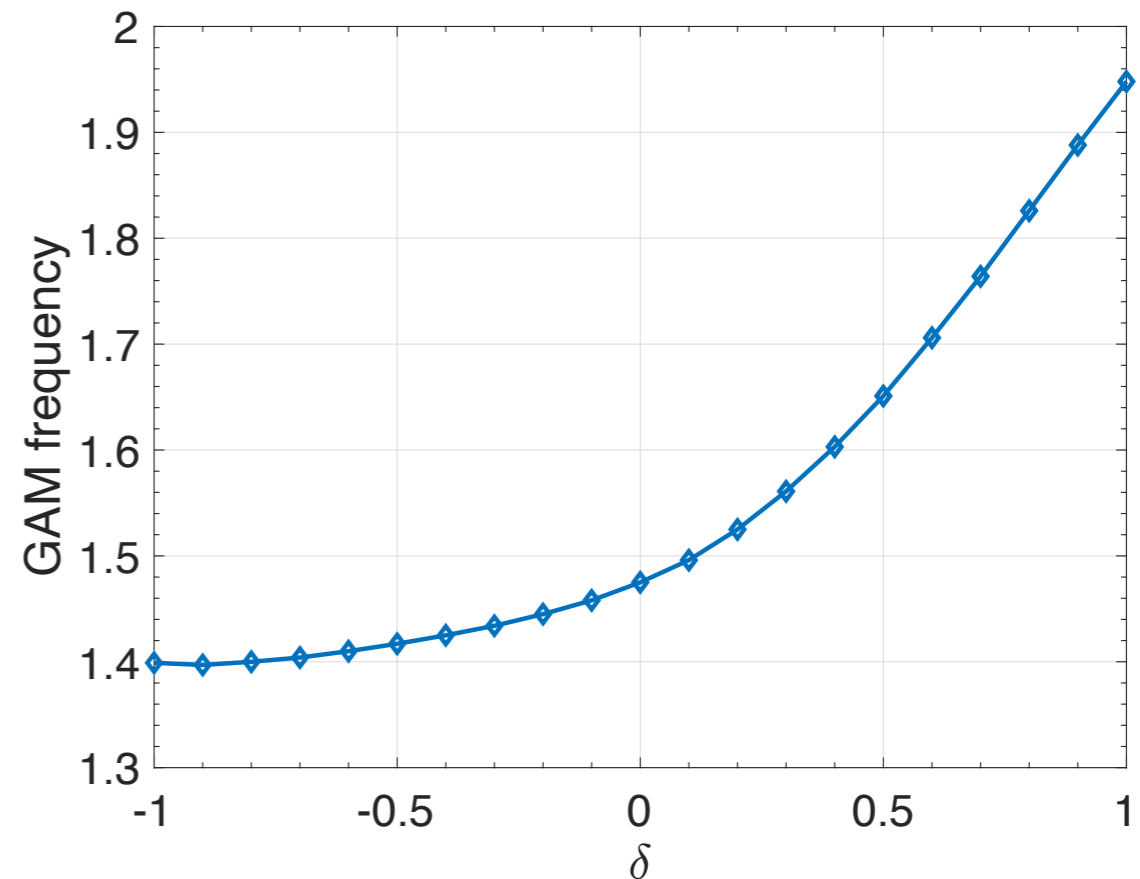
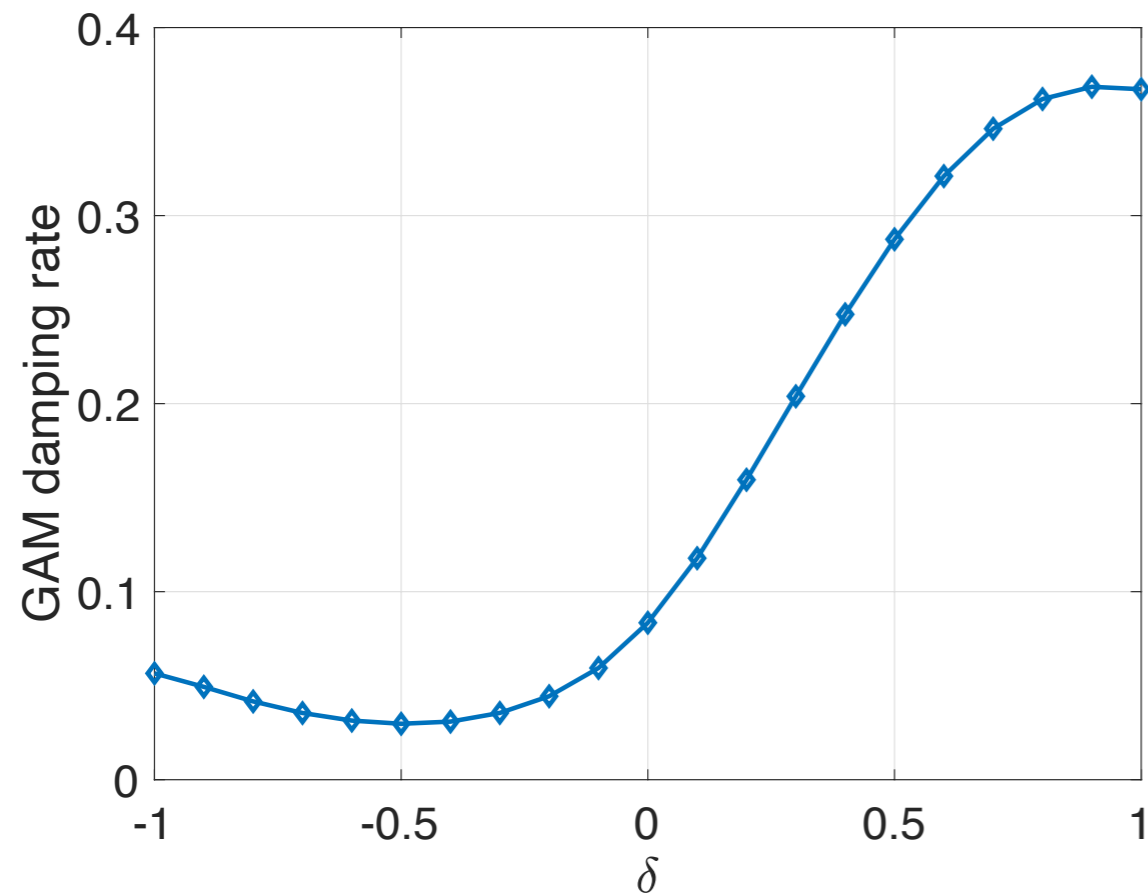
Zonal flows are reduced in NT

[Singh and Diamond NF 2022]



Wider magnetic well \rightarrow increased trapped fraction \rightarrow Stronger neoclassical screening \rightarrow reduced zonal flows

GAM frequency and damping rates reduced in NT



- GAM Landau damping is more strongly (~ 7 times) reduced than the GAM frequency for NT!
 \implies More coherent and stronger GAM ExB shearing field for NT than for PT!
 \implies NT plasma turbulence is likely saturated by GAMs!

...more work in progress...

Geometry dependence of mean ExB shearing rate ω_E

[Singh and Diamond NF 2023, under review]

ExB shearing rate in general axisymmetric toroidal geometry obtained from a 2-point correlation calculation: [Hahm & Burrell PoP 1995]

$$\omega_E = \left(\frac{\Delta\psi_0}{\Delta\zeta} \right) \frac{\partial^2}{\partial\psi^2} \Phi_0(\psi),$$

$\psi :=$ poloidal flux $\zeta :=$ toroidal angle

$\Phi_0 :=$ Mean electrostatic potential

$\Delta\psi_0 :=$ Turbulence correlation length in ψ

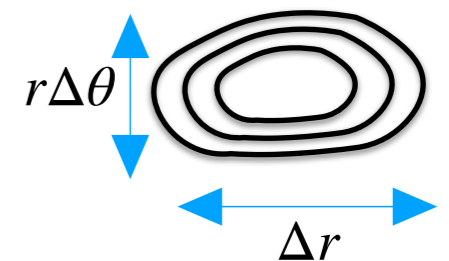
$\Delta\zeta :=$ Turbulence correlation in toroidal angle ζ

- $\frac{\partial^2}{\partial\psi^2} \Phi_0(\psi)$ is set by the radial force balance of ions - as usual!
- $\Delta\psi$ is related to turbulence radial correlation length Δr : $\Delta\psi = \Delta r \frac{\partial\psi}{\partial r}$, where ψ' is obtained from the definition of global safety factor q : $\psi' = \frac{I(\psi)}{2\pi q(\psi)} \oint d\theta \frac{\mathcal{J}}{R^2}$
- $\Delta\zeta$ is related to poloidal correlation angle $\Delta\zeta = \nu \Delta\theta$, where the local safety factor $\nu = \frac{I\mathcal{J}}{R^2\psi'}$

$$\text{Thus, } \omega_E = \frac{\Delta r}{\Delta\theta} \frac{R^2\psi'^2}{I\mathcal{J}} \frac{\partial^2}{\partial\psi^2} \Phi_0(\psi),$$

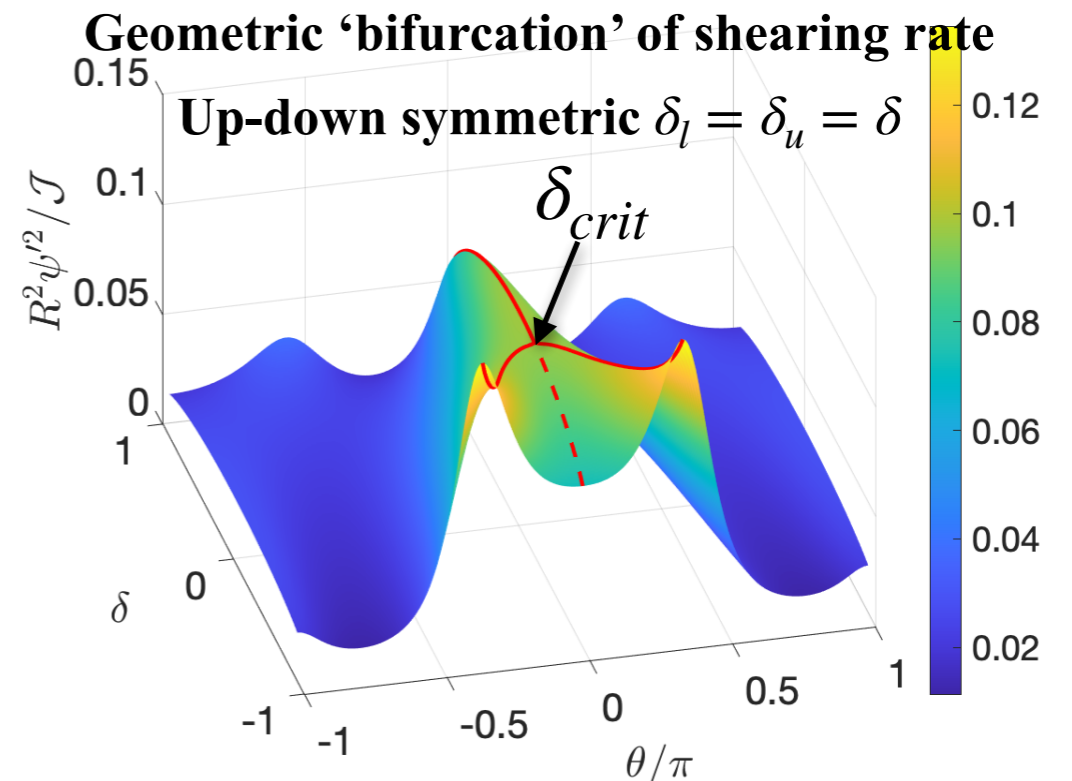
Geometry dependent factor

Calculated for Miller's equilibrium for **fixed** $\frac{\Delta r}{\Delta\theta}$ and $\frac{\partial^2}{\partial\psi^2} \Phi_0(\psi)$.

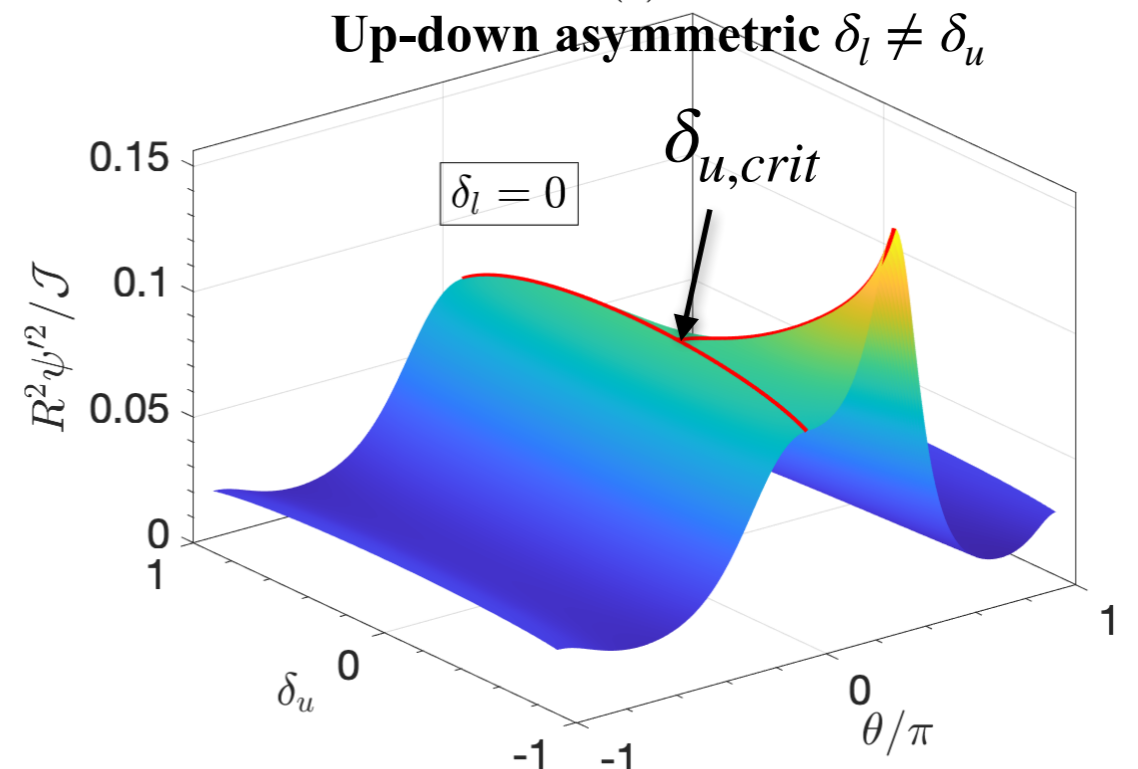
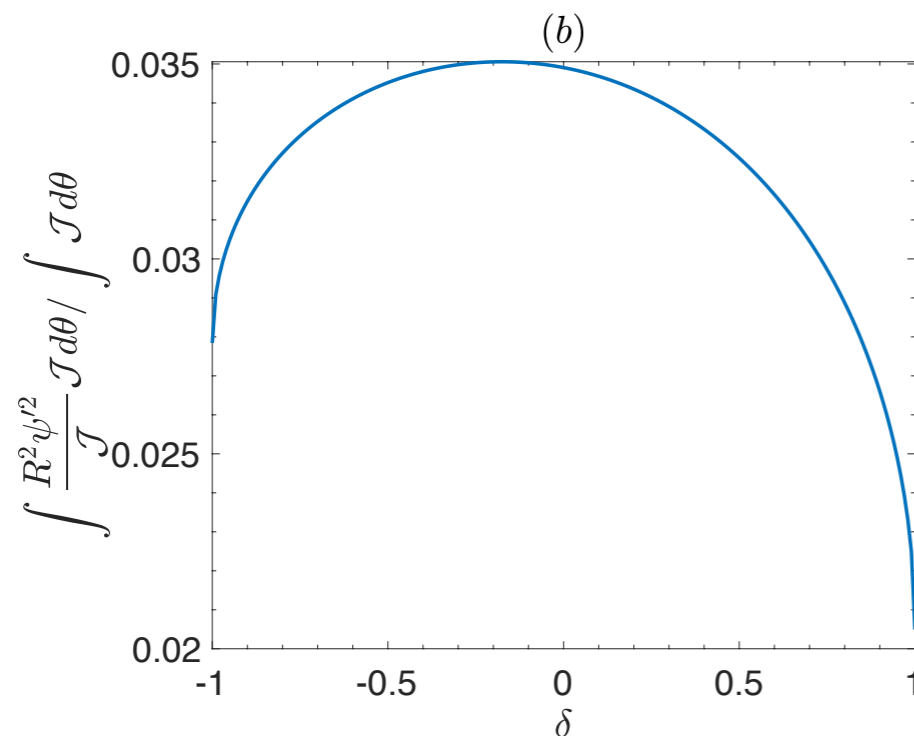


Variation of mean ExB shearing rate with triangularity δ

- ➔ • **Maximum shear off the outboard mid-plane for for $\delta < \delta_{crit}$ (\sim NT) \rightarrow Shearing is less effective for $k_x = 0$ modes i.e, the modes ballooning at $\theta = 0$.**
- **Shear at $\theta = 0$:**
 - \downarrow with increasing NT.
 - Weaker for NT than for PT. Note that fluctuations balloon at $\theta = 0$. Thus, shearing efficiency $\downarrow \implies P_{L \rightarrow H, th} \uparrow$ (!?).



(b)

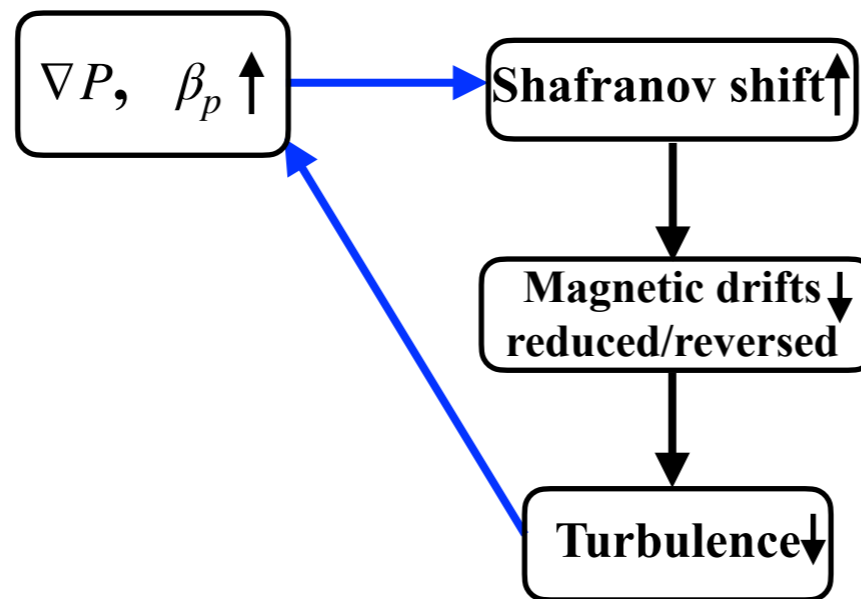


- Flux surface averaged shearing rate is higher for NT than for PT. -Global confinement ?!

- Maximum shear above the outboard mid-plane for for $\delta_u < \delta_{u,crit}$ & $\delta_l > \delta_u$

Shafranov shift induced transport bifurcation

- ITB formation in high- β_p regime is often linked to transport bifurcation due to turbulence stabilization by Shafranov shift due to mag drift reduction/reversal, *ignoring* the mean ExB shear effect. [Mike Beer *et al* PoP 1997, S Ding *et al* PoP 2017, J McClenaghan *et al* PoP 2019, G M Staebler *et al* PoP 2017]



Feedback loop for Shafranov shift induced transport bifurcation

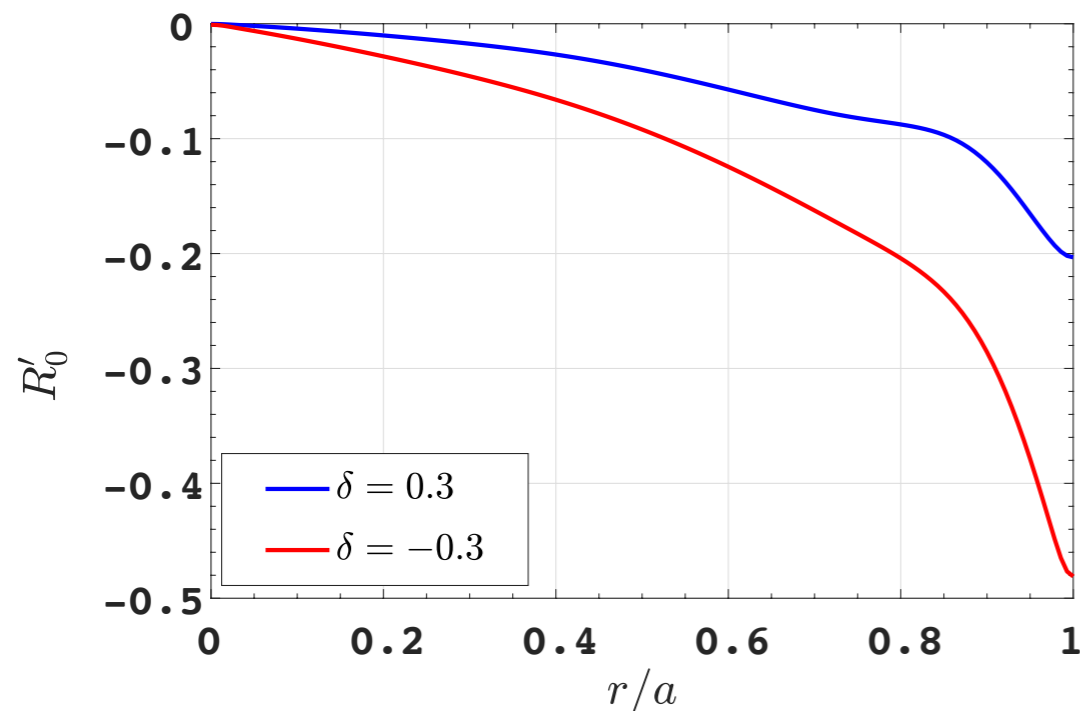
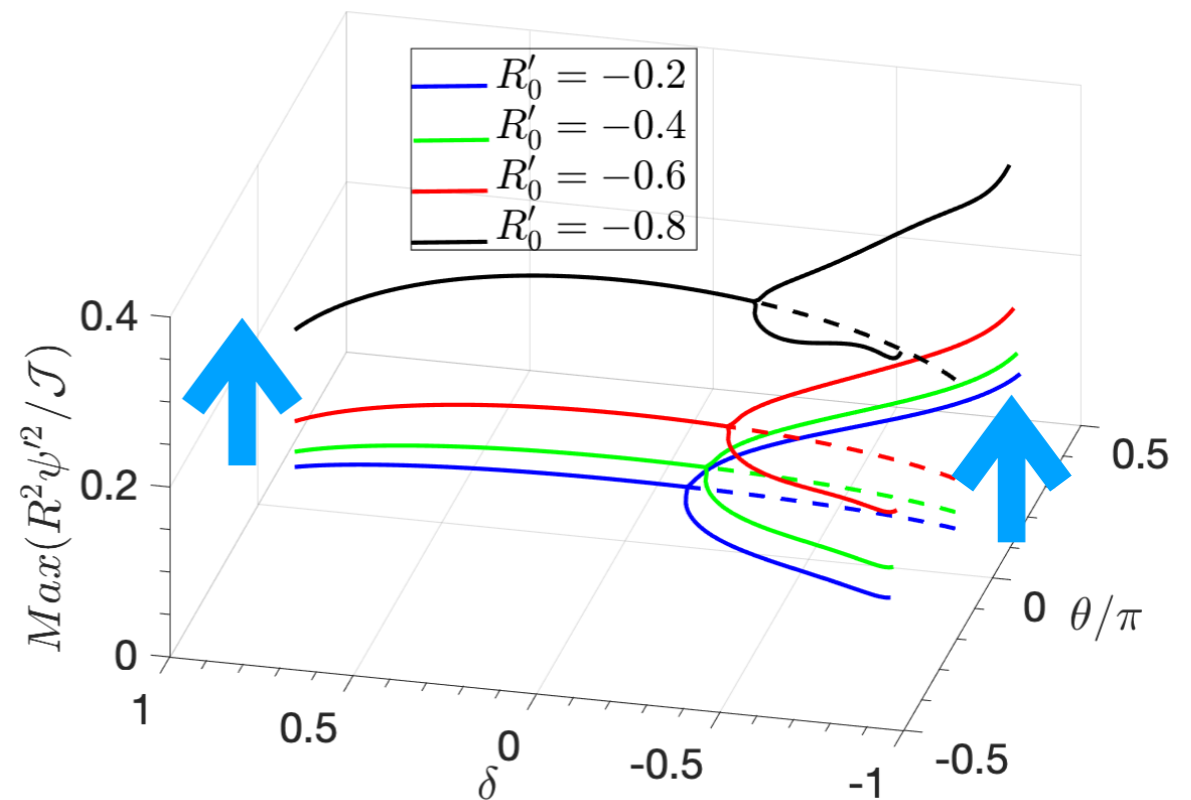
- But... like it or not - mean shear *exist* in high- β_p discharges!
- So **how does mean shear and Shafranov shift interact ?**
- **Interplay of mean ExB shear, Shafranov shift and NT?**

Variation of mean ExB shearing rate with Shafranov shift gradient R'_0

On increasing $-R'_0$:

- Shearing rate increases for all δ .
- δ_c moves toward δ^- on increasing $-R'_0$.

→ • Key reason → flux compression.



Shafranov shift gradient obtained using CHEASE code

Significant for:

- high β_p regime (i.e, RS ITB) as $R'_0 \propto \frac{r}{R_0} \beta_p$
- NT shapes
 - as $\beta_p(\delta^-) > \beta_p(\delta^+)$
 - Numerical MHD equilibrium study shows $R'_0(\delta^-) > R'_0(\delta^+)$ even for fixed β_p .

→ Even more significant for future NT+ITB discharges

Implications of Shafranov shift effect on ExB shear

- Shafranov shift affects turbulence in 2 distinct ways:

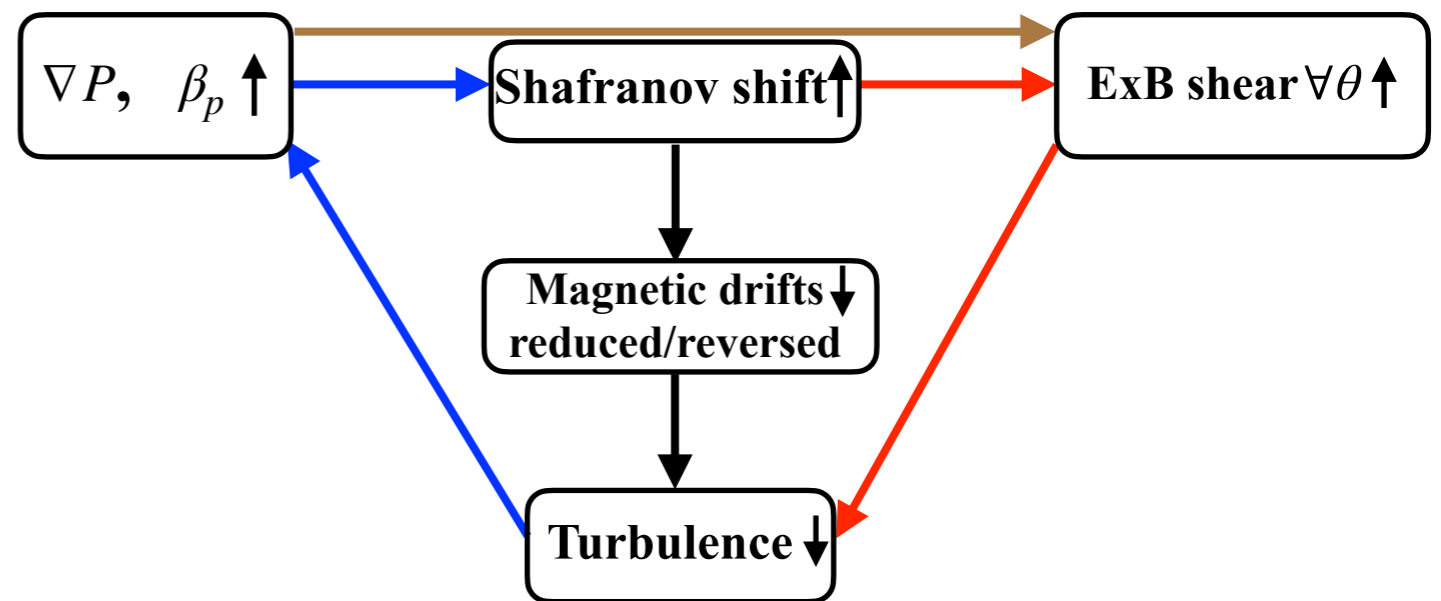
(I) Stabilizes turbulence by reduction/reversal of magnetic drifts

(II) Directly enhances the mean shear, \rightarrow additional turbulence suppression

Both can cause bifurcation to enhanced confinement state independently. Bifurcation by (I) is often invoked as a mechanism of confinement improvement in high- β_p regime, *ignoring* the mean shear effect.

Enhanced mean ExB shearing by Shafranov shift provides a +ve feedback on the feedback loop of the Shafranov shift induced transport bifurcation.

Shafranov shift also has a +ve effect on the mean ExB shear induced transport bifurcation, not only through a reduction of the linear growth rate but also through the enhanced ExB shearing rate.



Both (I) and (II) can work in tandem to reduce the ∇P_{crit} for the onset of ITB in reversed shear PT shape

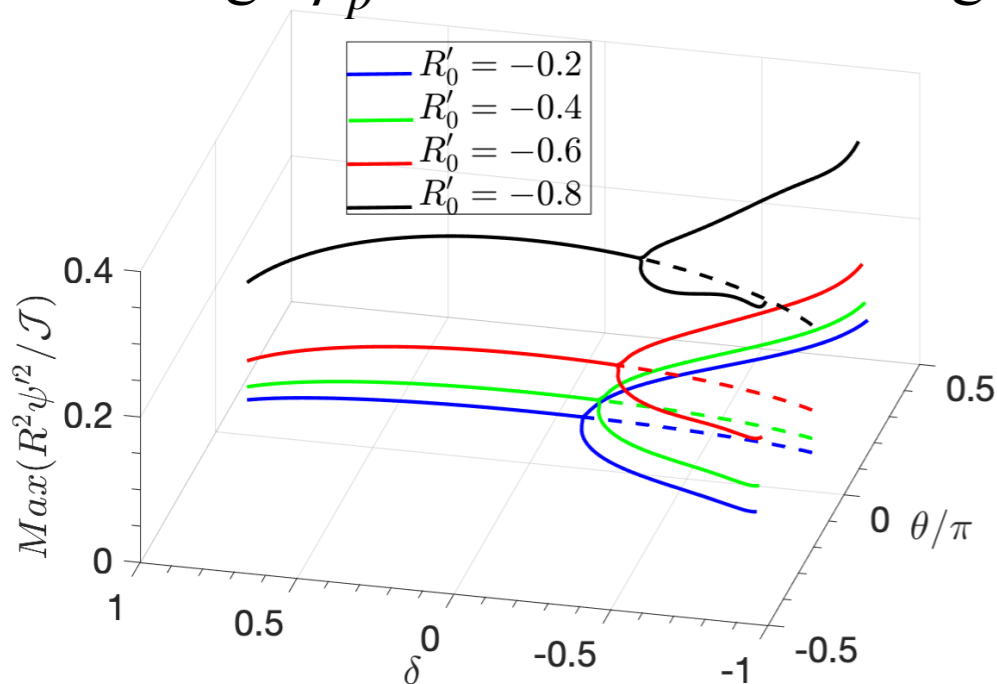
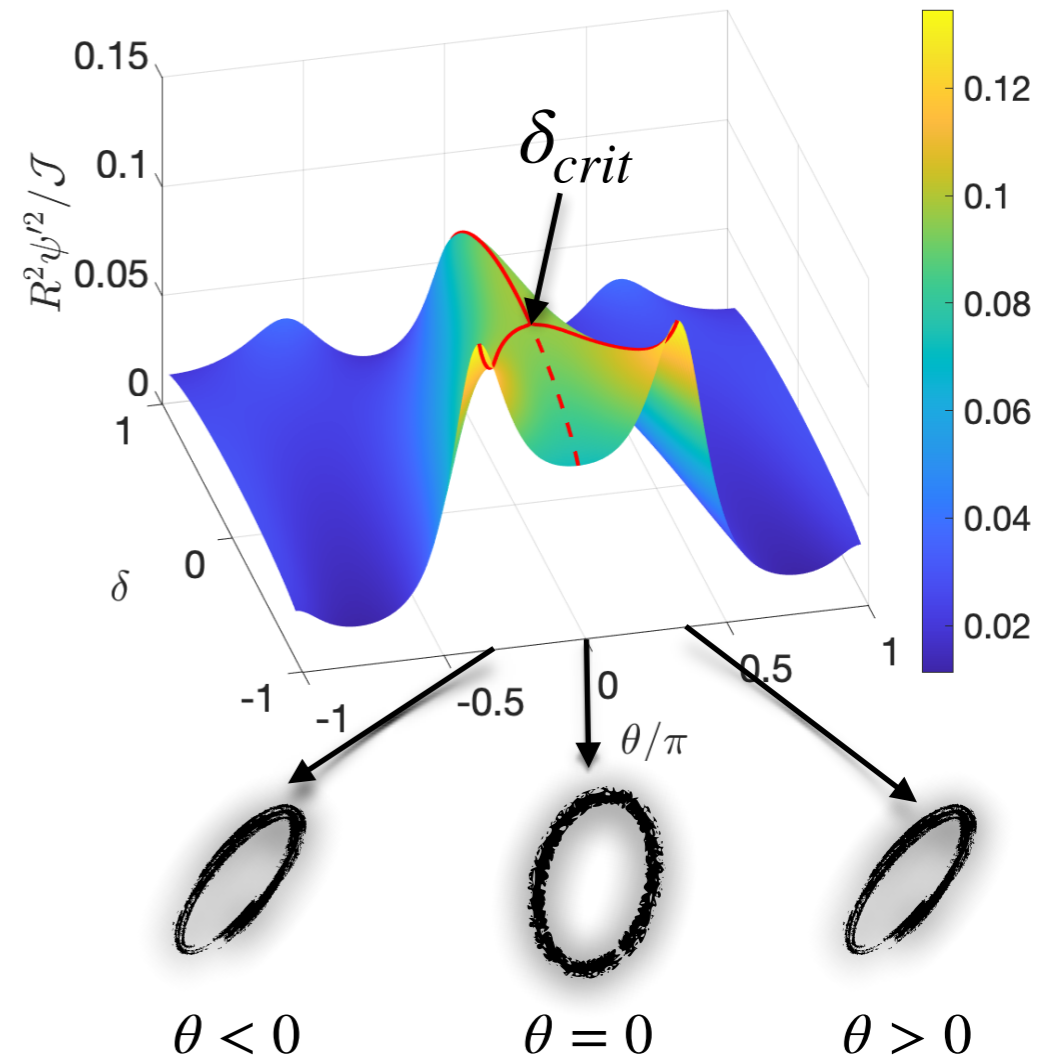
Conclusions

- **Zonal flows are weaker in NT than in PT** due to increased neoclassical screening, from an increase in trapped fraction in NT.
- **GAM frequency and Landau damping rates are significantly reduced in NT** due to reduction of both magnetic drift frequency and parallel transit frequency.
 - GAM is likely the dominant player for turbulence regulation in NT.
- **Maximum shear off the outboard mid-plane for $\delta < \delta_{crit} (\leq 0)$**
 - Up-down symmetry: Max shear located symmetrically above and below the outboard mid-plane for $\delta_u = \delta_l = \delta < \delta_{crit}$
 - Up-down asymmetry: Max shear located above the outboard mid-plane for $\delta_u < \delta_{crit}$ & $\delta_l > \delta_u$. Max shear located below the outboard mid-plane for $\delta_l < \delta_{crit}$ & $\delta_u > \delta_l$.
 - **Shearing is more effective for $k_x \neq 0$ modes for NT. Are these relevant?**
 - Shear at $\theta = 0$ decreases with increasing NT. Fluctuations balloon at $\theta = 0$. Thus, **shearing efficiency $\downarrow \implies P_{L \rightarrow H, th} \uparrow (!?)$. Is this sufficient?**
- **Direct effect of Shafranov shift gradient $-R'_0$ on shearing rate: Shearing rate increases with increasing $-R'_0$ for all δ .** Key reason \rightarrow flux compression. Significant for high β_p regime and NT shapes.

These results has implications not just for confinement & L-H transition for NT but also for ITB discharges in PT and NT(proposed), and NT core and and pedestal.

For the experimentalists

- **Mean ExB Shearing is maximal off the mid-plane for NT:** \implies Eddy tilting should be strongest off the mid-plane.
 - Direct imaging using gas-puffing.
 - Joint pdf of radial and poloidal velocity fluctuations (i.e., \tilde{v}_r & \tilde{v}_θ) should show max tilting (most-correlated) off the mid-plane for NT.
 - Up-down asymmetric tilting distribution for $\delta_u \neq \delta_l$
- **Shafranov shift gradient R'_0 directly boosts the mean ExB shear:**
 - Re-assess the role of mean ExB shear in high- β_p reverse shear discharges.



- **Variation of GAM energy to ZF energy as PT \rightarrow NT, using fluctuation diagnostics.**
- **Radial correlation length of ZF vs δ , frequency resolved Reynolds power vs δ , using BES diagnostics.**